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FT-Raman, FT-IR spectra and total energy distribution of 3-pentyl-2,6-diphenylpiperidin-4-one: DFT method

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FT-Raman and FT-IR spectra were recorded for 3-pentyl-2,6-diphenylpiperidin-4-one (PDPO) sample in solid state. The equilibrium geometries, harmonic vibrational frequencies, infrared and the Raman scattering intensities were computed using DFT/6-31G(d,p) level. Results obtained at this level of theory were used for a detailed interpretation of the infrared and Raman spectra, based on the total energy distribution (TED) of the normal modes. Molecular parameters such as bond lengths, bond angles and dihedral angles were calculated and compared with X-ray diffraction data. This comparison was good agreement. The intra-molecular charge transfer was calculated by means of natural bond orbital analysis (NBO). Hyperconjugative interaction energy was more during the π – π^* transition. Energy gap of the molecule was found using HOMO and LUMO calculation, hence the less band gap, which seems to be more stable. Atomic charges of the carbon, nitrogen and oxygen were calculated using same level of calculation.

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1. Introduction

Piperidones exhibit a wide spectrum of biological activities and form an essential part of the molecular structures of important drugs. Molecular geometry critically influences biological activity. Attention has been focused on structure-activity relationships. Piperidines with crowded groups at C_3 and C_5 have enhanced biological activity compared to other piperidines [\[1\].](#page-9-0) 2,6-Disubstituted piperidin-4-ones are regarded as an important framework and served as precursors for chiral biologically active natural alkaloids [\[2\].](#page-9-0) The biological activities of piperidones were found to be excellent if 2- and/or 6-positions are occupied by aryl groups [\[3\].](#page-9-0) Accordingly, anti-bacterial and anti-fungal activities of 2,6-diarylpiperidin-4-ones and their derivatives have been explored well [\[4,5\].](#page-9-0) Stereochemistry of N-benzoyl-2r,6c-diphenylpiperidin-4-one oxime, N-benzoyl-3t-methyl-2r,6c-diphenylpiperidin-4-one oxime, N-benzoyl-3tethyl-2r,6c-diphenylpiperidin-4-one oxime (3), N-acetyl-2r,6cdiphenylpiperidin-4-one oxime and N-acetyl-3t-methyl-2r,6cdiphenylpiperidin-4-one oxime have been studied using 1 H, 13 C and two-dimensional NMR spectra [\[6\].](#page-9-0) The main goal of this work is to record, simulate and interpret the vibrational spectra for the title compound, which has not been presented before. We also wanted to shed a light on the crystal and vibrational spectral data (FT-Raman and FT-IR) with the results of theoretical calculations.

2. Experimental details

2.1. Synthesis of 3-pentyl-2,6-diphenylpiperidin 4-one (PDPO) [\[7\]](#page-9-0)

A mixture of ammonium acetate (3.85 g, 0.05 mol), benzaldehyde (10.6 ml, 0.1 mol) and 2-octanone (6.4 ml, 0.05 mol) in distilled ethanol was heated to boiling. After cooling the viscous liquid obtained was dissolved in diethyl ether (200 ml) and was shaken with 2 ml concentrated hydrochloric acid. The precipitated hydrochloride of the title compound was removed by filtration and washed with 40 ml mixture of ethanol and diethyl ether (1:1) and then with diethyl ether to remove most of the coloured impurities [\[7\].](#page-9-0) The base was liberated from an alcoholic solution by adding aqueous ammonia and then diluted with water. It was purified by column chromatography, using a n-hexane-ethyl acetate mixture as the solvent. The yield of the compound was 80%.

2.2. FT-Raman and FT-IR measurement

The FT-Raman spectrum of PDPO was recorded using the 1064 nmline of Nd:YAG laser as excitation wavelength in the region 10–3500 cm−¹ on a Bruker model RFS100/S spectrophotometer

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Fig. 1. (a) Relative energy–dihedral angle curve in C₃₆–C₃₉ bond. (b) Relative energy–dihedral angle curve in C₅–C₃₅ bond. (c) Relative energy–dihedral angle curve in C₄–C₁₃ bond. (d) Relative energy–dihedral angle curve in C₃–C₂₄ bond. (e) Relative energy–dihedral angle curve in C₃₉–C₄₂ bond. (f) Relative energy–dihedral angle curve in C₄₂–C₄₇ bond.

equipped with FRA 106 FT-Raman module accessory. The spectral measurements were carried out at Sree Chitra Tirunal Institute for Medical Sciences and Technology, Poojappura, Thiruvananthapuram, Kerala, India. The FT-IR spectrum of this compound was recorded in the region 400–4000 cm−¹ on an IFS 66V spectrophotometer using the KBr pellet technique. The spectrum was recorded at room temperature, with a scanning speed of 10 cm^{-1} per minute and at the spectral resolution of 2.0 cm−¹ in CISL Laboratory, Annamalai University, Tamil Nadu, India.

3. Computational details

The entire calculations were performed at DFT levels on a Pentium 1V/3.02 GHz personal computer using Gaussian 03W [\[8\]](#page-9-0) program package, invoking gradient geometry optimization [\[8,9\].](#page-9-0) Initial geometry generated from standard geometrical parameters was minimized without any constraint in the potential energy surface at DFT level, adopting the standard 6-31G(d,p) basis set. The optimized structural parameters were used in the vibrational frequency calculations at the DFT level to characterize all stationary points as minima. Then, vibrationaly averaged nuclear positions of PDPO were used for harmonic vibrational frequency calculations resulting in IR and Raman frequencies together with intensities and Raman depolarization ratios. In this study, the DFT method (B3LYP) was used for the computation of molecular structure, vibrational frequencies and energies of optimized structures. The vibrational modes were assigned on the basis of TED analysis using som program [\[10\].](#page-9-0)

Fig. 2. Optimized molecular structure of 3-pentyl-2, 6-diphenylpiperidin-4-one.

It should be noted that Gaussian 03W package able to calculate the Raman activity. The Raman activities were transformed into Raman intensities using Raint program [\[11\]](#page-9-0) by the expression:

$$
I_i = 10^{-12} \frac{(v_0 - v_i)^4}{v_i \cdot S} \tag{1}
$$

where I_i is the Raman intensity, S is the Raman scattering activities, v_i is the wavenumber of the normal modes and $v_{\mathbf{0}}$ denotes the wavenumber of the excitation laser [\[12\].](#page-9-0)

4. Results and discussion

4.1. Conformational analysis

The chair conformer of piperidine molecule is the most stable conformer. Therefore, we neglected other conformations that differ from the chair (boat, envelope or twist boat) because of their high energy. Moreover, it has two possible chair conformations, which differ in the axial (A) or equatorial (E) positions of the N–H group [\[13–15\].](#page-9-0) Piperidine molecules show the equatorial form of NH of chair conformer as the most stable. Piperidine molecule adopts the NH equatorial position of the chair conformer. Then, in order to reveal all possible conformations of studied molecule, a detailed potential energy surface (PES) scan in six dihedral angles was performed. This scan was carried out by relaxed PES scanning calculations in all geometrical parameters by changing the torsion angle for every 10◦ rotation around the bond. The shape of the potential energy as a function of the dihedral angle is illustrated in [Fig.](#page-1-0) 1a–f. The curves between relative energy and dihedral angles $(a \rightarrow H_{11}-C_3-C_{24}-C_{25}$ and $c \rightarrow H_9-C_4-C_{13}-C_{14}$) are shown in [Fig.](#page-1-0) 1a and c. As seen in [Fig.](#page-1-0) 1a and c, $H_{11}-C_3-C_{24}-C_{25}$ dihedral angle of phenyl ring is attached with C_3 is determined at 160 \textdegree for B3LYP level of theory. H₉–C₄–C₁₃–C₁₄ dihedral angle is predicted at 10 $^{\circ}$. In the optimized structure, $H_{11}-C_3-C_{24}-C_{25}$ and $H_9-C_4-C_{13}-C_{14}$ dihedral angles are predicted at 162.24◦ and 9.50◦, respectively.

4.2. Molecular geometry

The optimized geometrical parameters and structure of PDPO was calculated at 6-31G(d,p) level, are given in [Table](#page-3-0) 1 and Fig. 2, respectively. Geometrical parameters such as bond lengths, bond angles and dihedral angles are also given along with its single crystal X-ray diffraction data. The bond length of C_1 -O₅₀ is about 1.218 in B3LYP method. And its corresponding experimental value

Fig. 3. Theoretical (a) and experimental (b) FT-IR spectrum of PDPO.

is 1.214Å [\[7\].](#page-9-0) Similarly the C_3-N_{12} and C_4-N_{12} bond distances are calculated at 1.465 and 1.470 Å (DFT), which are in agreement with X-ray data. The bond distance of C–C is usually observed as ~1.400Å. In the present investigation, bond lengths of C_1-C_2 , C_1 – C_5 , C_2 – C_3 , C_3 – C_{24} , C_4 – C_5 , C_4 – C_{13} and C_5 – C_{35} are in line with literature values. On the other hand, the bond distances $(C_{13}-C_{14},$ $C_{13}-C_{15}$, $C_{14}-C_{16}$, $C_{15}-C_{18}$, $C_{16}-C_{20}$, $C_{18}-C_{20}$, $C_{25}-C_{27}$, $C_{26}-C_{29}$, $C_{27}-C_{31}$ and $C_{29}-C_{31}$) of the sixmembered rings are approximately 1.39 \AA /B3LYP with few exceptions. These values are in agreement with literature values [\[7\].](#page-9-0) Crystal data [\[7\]](#page-9-0) reveal that the C–H bond distances are \sim 1.00Å, which is supported by the calculated values. The calculated angles 121.85◦ and 122.91◦ (DFT), are belongs to C_2 – C_1 – O_{50} and C_5 – C_1 – O_{50} , respectively. And their corresponding literature values are 121.93◦ and 122.01◦. These larger bond angles are may be due to electron density in oxygen atom. The bond angle $C_{14}-C_{13}-C_{15}$: 118.705/B3LYP is less when comparing with other bond angles $C_{13}-C_{14}-C_{16}$, $C_{13}-C_{15}-C_{18}$, $C_{14}-C_{16}-C_{20}$, C₁₅–C₁₈–C₂₀ and C₁₆–C₂₀–C₁₈ (∼120°). It may be due to the phenyl ring is attached with C_4 . Similar trend has been observed in the second phenyl ring. The bond angles of C–C–H are ∼109◦ except in both phenyl rings (∼120◦) which are in agreement with literature values [\[7\].](#page-9-0) The dihedral angles of title molecule were calculated, and some of them were compared with available X-ray diffraction data as shown in [Table](#page-3-0) 1.

4.3. Vibrational assignments

Synthesized PDPO, consists 51 atoms and hence 147 normal modes of vibrations and the molecule belongs to C_1 symmetry. The fundamental vibrational wavenumbers of PDPO was calculated by $DFT(B3LYP/6-31G(d,p))$ is given in [Table](#page-5-0) 2. The resulting vibrational wavenumbers for the optimized geometries, IR intensities as well as Raman scattering activities and experimental FT-IR, FT-Raman frequencies are also listed. Experimental and theoretical spectra of title compound have been shown in Fig. 3 (FT-IR) and [Fig.](#page-4-0) 4 (FT-Raman). The normal modes of vibration were assigned on the basis of TED. To bring the theoretical values closer to experimental values, we used the scale factor: 0.9608.

4.3.1. N–H vibrations

The N–H stretching vibration [\[13,14\]](#page-9-0) appears strongly and broadly in the region 3500–3300 cm−1.Y. Erdogdu et al., assigned v_{N-H} mode in the region 3500–3300 cm⁻¹ [15]. In this study, the frequency was observed as weak and narrow band in both FT-IR and FT-Raman, where the frequencies are attributed to 3316 cm⁻¹ and

Table 1 Bond lengths, bond angles and dihedral angles of PDPO.

Table 1 (Continued)

Parameters Bond length (\AA)	Exp ^a	$B3LYP/6-31G(d,p)$	Parameters Angles Contd.	Exp ^a	$B3LYP/6-31G(d,p)$
C_{25} -C ₂₇ -C ₃₁	120.8	120.0	$N_{12}-C_{4}-C_{5}-C_{1}$	-51.99	-54.54
C_{26} -C ₂₉ -C ₃₁	120.7	120.2	N_{12} – C_4 – C_5 – C_{35}	-178.04	-179.93
$C_5 - C_{35} - C_{36}$	113.4	113.6	$C_{13}-C_{4}-C_{5}-C_{1}$	-172.73	-175.84
$C_{35}-C_{36}-C_{39}$	113.7	112.9	$C_{13}-C_{4}-C_{5}-C_{35}$	61.21	58.77
$C_{36} - C_{39} - C_{42}$	112.8	113.4	$C_{13} - C_{4} - N_{12} - C_{3}$	-171.28	-172.35
$C_{39} - C_{42} - C_{47}$	113.7	113.2			
$C_2 - C_1 - O_{50}$	121.9	121.8			
$C_2 - C_1 - O_{50}$	122.0	122.9			

^a Ref. [\[7\].](#page-9-0)

Fig. 4. Theoretical (a) and experimental (b) FT-Raman spectrum of PDPO.

3317 cm⁻¹ respectively. The corresponding theoretical frequency for v_{N-H} mode is about 3384 cm⁻¹, which shows positive deviation of [∼]⁶⁸ cm−¹ from the experimental value. The scissoring mode of N–C–H is appeared at 1407 cm⁻¹ (FT-IR-strong) and 1409 cm⁻¹ (FT-Raman-weak), while the harmonic scissoring vibration existed at 1411 cm⁻¹ (mode no. 104). These δ_{C-N-H} wavenumbers are also find support from the literature. Out-of-plane bending modes (γ_{C-N-H}) are calculated at about 765 and 771 cm−¹ (mode nos. 44, 45), these vibrations are in line with the observed FT-IR (765 cm⁻¹) and FT-Raman (785 cm⁻¹) bands.

4.3.2. Methyl and methylene group vibrations

Methyl groups are generally referred to as electron donating substituents in the aromatic ring system [\[16\].](#page-9-0) In acetates, the asymmetric vibrations of the methyl group are expected to occur in the region 2940–3040 cm−¹ and symmetric vibrations are in the region 2910–2930 cm⁻¹, and usually the bands are weak [\[17\].](#page-9-0) Aromatic acetyl substituents absorb in a narrow range 3000–3020 cm⁻¹ absorption sometimes coincides with a CH stretching mode of the ring [\[17\].](#page-9-0) The title molecule possesses methyl (CH_3) and methylene (CH2) groups. Methyl group symmetric stretching vibrations are appeared at 2914 cm−¹ as a strong intense band in FT-IR and 2913 cm−¹ as a very weak band in FT-Raman spectrum. While the scaled harmonic frequency 2917 cm−¹ (mode no: 127) with considerable intensity is in line with experimental value. The asymmetric of CH₃ harmonic frequency 2987 cm⁻¹ (mode no: 135) is coincide well with experimental values (FT-IR: 2989/FT-Raman: 2990 cm⁻¹).

The asymmetric and symmetric $CH₂$ stretching vibrations are normally appear in the region 3100–2900 cm⁻¹ [\[18\].](#page-9-0) According to the literature [\[18\],](#page-9-0) the observed bands 2870 cm−¹ (weak), 2924 cm⁻¹ (strong) in FT-IR and 2890 cm⁻¹ (weak) FT-Raman are in agreement with the theoretical values in the range of 2891–2937 cm−¹ (mode nos: 123–126, 128–130). The FT-IR band 2956 cm−¹ (strong) and its corresponding theoretical value 2960 cm⁻¹ (mode no: 132) are attributed v_{asy} CH₂ vibration. The TED value shows that these vibrations are pure. In aromatic compounds the v_{C-H} , β_{C-H} and γ_{C-H} modes are appeared in the range of 3000–3100 cm⁻¹, 1000–1300 and 750–1000 cm⁻¹, respectively [\[19–21\].](#page-9-0) The C–H stretching vibrations appeared at 2810, 3048–3077 cm−¹ (mode nos: 122, 137–144). The observed frequencies 2855, 3061 cm⁻¹ (FT-IR) and 2856, 3042, 3057 cm⁻¹ (FT-Raman) are belongs to C–H stretching mode. The C–H in-plane bending vibrations appeared in the range 1123–1170 cm⁻¹ (mode nos: 76, 79, 81) and their corresponding experimental wavenumbers 1118, 1145, 1155 (FT-IR) and 1129, 1156 cm−¹ (FT-Raman) are in consistent with computed values. The assignments also find support from the literature [\[22,23\].](#page-9-0)

The scissoring mode of the $CH₂$ group gives rise to a characteristic band near 1415 cm^{-1} in IR and 1400 cm^{-1} in Raman spectra. The twisting, wagging and rocking vibrations appear in the region 1400–900 cm−¹ [\[24\].](#page-9-0) The broadening and intensity decreases were observed for the bands at 1470 and 1450 cm−¹ corresponding to CH₂ scissoring modes [\[22\].](#page-9-0) In the present investigation, δ _{CH2} mode appear at 1436 and 1456 cm−¹ as medium band at FT-IR, and their FT-Raman counter parts are 1452 and 1432 cm−¹ as weak bands. These experimentalfrequencies are in agreement with mode nos: 106 and 111 of B3LYP. A major coincidence of theoretical values with that of experimental observation is found for ω_{CH2} (1348:B3LYP-mode no: 100/1347 cm−1: FT-IR). These assignments find support from the literature $[22]$. In the case of CH₂ twisting mode the vibrational frequencies observed at 1274, 1303 cm−¹ (FT-IR) and 1304 cm−¹ (Raman) are in agreement with the calculated frequencies in the range of 1278–1300 cm−¹ (mode no: 90, 92–94/DFT). The out-of-plane bending mode of C–H always lies in the lower side of spectra. In the present study, the harmonic wavenumbers (mode nos: 60–58, 56, 54–52, 48, 47, 43–40) in the range of 962–706 cm⁻¹ are assigned to $γ_{C-H}$ mode, which find support from the observed FT-IR frequencies: 912, 889, 757 and 738 cm−1.

4.3.3. $C=0$. $C-N$. $C=C$ vibrations

Stretching vibration of carbonyl group $C=O$ can be observed as a very strong band in both FT-IR and FT-Raman spectra at 1665 cm⁻¹ [\[22\].](#page-9-0) The carbonyl stretching C=O vibration [\[17,25\]](#page-9-0) is expected to occur in the region 1715–1680 cm−1. The deviation of the calculated wavenumbers for this mode can be attributed to the underestimation of the large degree of π -electron delocalization due to conjugation of the molecule [\[26\].](#page-9-0) The literature

Table 2 (Continued)

Table 2 (Continued)

v: stretching, δ : bending, T: torsion, vw: very week, w: week, m: medium, s: strong, vs: very strong.

^a Scaling factor: 0.9608 [\[28\].](#page-9-0)

b Relative absorption intensities normalized with highest peak absorption equal to 100.

 $\,^{\rm c}$ Relative Raman intensities calculated by Eq. [\(1\)](#page-2-0) and normalized to 100.

 d Total energy distribution calculated B3LYP 6-31G(d,p) level, TED less than 10% are not shown.

Second order perturbation theory analysis of Fock matrix in NBO basis (PDPO).

 $E^{(2)}$ means energy of hyperconjucative interactions (stabilization energy).

 b Energy difference between donor and acceptor i and j NBO orbitals.

 ϵ $F(i, j)$ is the Fock matrix element between *i* and *j* NBO orbitals.

[\[27\]](#page-9-0) reveals that the normal esters are characterized by the strong IR absorption due to the $C=O$ stretching vibration in the range of 1750–1735 cm−1. In this study, we have observed stretching of C=O at 1715 cm⁻¹ as very strong in FT-IR and 1714 cm⁻¹ as medium intense band in FT-Raman, while the computed frequency is 1732 cm−¹ (mode no: 120) and its TED value (90%). The C–C stretching in phenyl ring and methylene chain is calculated in the range of 1035–864 cm−¹ (mode nos: 69–64, 62, 61, 57, 55, and 52–50). These vibrations are inline withexperimental values (1028, 1001, 986, 973, 889, 857: FT-IR and 1030, 1003, 988, 919 cm−1: FT-Raman) and also in consistent with literature values [\[22,23\].](#page-9-0) These assignments are further supported by the TED values.

The identification of C–N vibration is a very difficult task, since mixing of several bands are possible in this region. However, with the help of theoretical calculation (DFT), the C–N stretching vibrations are calculated. The C–N stretching vibration coupled with scissoring of N–H, is moderately to strongly active in the region 1275 ± 55 cm⁻¹ [\[17\].](#page-9-0) In the present investigation C–N stretching frequencies are observed at 1046, 1056, 1068,1086 and 1110 cm−¹ by FT-IR and their corresponding calculated wavenumbers appeared in the range of 1049–1104 cm−¹ (mode nos: 70–76). These experimental values of C–N stretching mode show good agreement with theoretical values. The ν C–N stretching vibration normally appears around 1300 cm⁻¹ [17]. In this work the vC-N frequencies are moderately lowered, which may be due to the mass effect around nitrogen atom.

5. NBO analysis

The hyperconjugation may be given as stabilizing effect that arises from an overlap between an occupied orbital with another neighboring electron deficient orbital when these orbitals are properly orientation. This non-covalent bonding–antibonding interaction can be quantitatively described in terms of the NBO analysis, which is expressed by means of the second-order

perturbation interaction energy $(E^{(2)})$ [\[29–32\].](#page-9-0) This energy represents the estimation of the off-diagonal NBO Fock matrix elements. It can be deduced from the second-order perturbation approach [\[33\]](#page-9-0)

$$
E^{(2)} = \Delta E_{ij} = q_i \frac{F(i,j)^2}{\varepsilon_j - \varepsilon_i} \tag{2}
$$

where q_i is the donor orbital occupancy, ε_i and ε_i are diagonal elements (orbital energies) and $F(i, j)$ is off diagonal NBO Fock matrix elements. In this present study we dealt with NBO analysis. Especially the amount of energy transfer from π bond orbital to anti bond $\bm{\pi}^*$ orbital, the stabilization energy $E^{(2)}$ associated with hyperconjugative interaction, $LPO(2) \rightarrow C_1-C_2$, and C_1-C_5 are obtained as 17.62 and 15.93 kJ/mol, respectively. The bond $C_{13}-C_{15}$ with electron density 1.656e, stabilize the energy of 19.77 and 20.59 kJ/mol to its acceptor anti bonding orbitals of C_{14} - C_{16} and C_{18} - C_{20} , respectively. These interactions are observed as an increase in electron density (ED) in C–C antibonding orbital that weaken their bonds [\[34\].](#page-9-0) This investigation clearly demonstrates that the occupancy value of bonding orbitals make sure the hyperconjugative interaction with maximum stabilization between filled and unfilled subsystem of the molecule. The ED of $C_{14}-C_{16}$ donor bond has \sim 1.669e, on the other hand its antibond ED (π^* C₁₃–C₁₅ and $C_{18}-C_{20}$) posses ~0.342 and 0.332e, and their $E^{(2)}$ energies are 20.99 and 19.69 kJ/mol respectively. From the NBO analysis, the lower the ED of donor with larger the ED of acceptor have maximum delocalization and become strong bond interaction. The higher the ED value with lower $E^{(2)}$ energy which becomes lesser interaction and hence it shifts the vibrational frequencies from the actual frequencies. It is evident that the C_3-N_{12} (1.979e) and C_4-N_{12} (1.981e) bond stretching vibration (in the range from 1049 to 1123; mode nos. 70–76) lowers from the normal C–N bond stretching (1300 cm⁻¹) [\[17\].](#page-9-0) This may be due to the lesser hyperconjugative interaction between C–N donor bonds to C–C acceptor bands. The $E^{(2)}$ values and types of the transition are shown in [Table](#page-7-0) 3.

6. HOMO–LUMO

The frontier molecular orbitals play an important role in the electric and optical properties, as well as in UV–vis spectra and chemical reactions [\[35\].](#page-9-0) The analysis of the wave function indicates that the electron absorption corresponds to the transition from the ground to the first excited state and is mainly described by one electron-excitation from the highest occupied molecular orbital (HOMO) to the lowest unoccupied orbital (LUMO) [\[36\].](#page-9-0) The energy gap for PDPO was calculated using B3LYP/6-31G(d,p) level. The bioactivity and chemical activity of the molecule depends on the eigen value of HOMO, LUMO and energy gap. LUMO as an electron acceptor represents the ability to obtain an electron; donor represents the ability to donate an electron. The frontier molecular orbitals are shown in Fig. 5. From the molecular orbital analysis the highest occupied level is 87 this locates over the C–N–C group. And the 88 is the excited frontier orbital (LUMO- π^*), this orbital located over the C₁–O₅₀ and carbon atoms in phenyl ring. The energy difference between the HOMO and LUMO is about 5.288 eV. The frontier molecular orbital of PDPO (HOMO–LUMO) is shown in Fig. 5.

HOMO energy $= -6.169$ eV

LUMO energy $= -0.881$ eV

Energy gap $= 5.288$ eV

The smaller band gap energy increases the stability of the molecule. The charge distribution of the molecule has calculated

LUMO (-0.881 eV)

Fig. 5. The frontier molecular orbital of PDPO (HOMO–LUMO).

Table 4 Atomic charge of PDPO.

using B3LYP/6-31G(d,p) level. This calculation depicts the charges of the every atom in molecule. Distribution of positive and negative charges is the cause, to increase or decrease of bond length. The atomic charges of carbon, nitrogen and oxygen are listed in Table 4, in which nitrogen atom has maximum negative charge of −0.515 and −0.465 a.u., for oxygen atom. The HOMO part is located over the $N_{12}-C_3$, $N_{12}-C_4$ orbital, is mainly due to the lone pair of electron. Some of the carbon atoms have only positive charge about C₁ (0.402), C₄ (0.022), C₁₃ (0.060) and C₂₄ (0.099 a.u.). This clearly explains that the LUMO exist in those areas. The Mulliken charge plot is shown in [Fig.](#page-9-0) 6.

7. Conclusion

All possible conformers are calculated by changing the torsion angle rotation with respect to bond. The calculated bond parameters are compared with reported X-ray diffraction data. All the vibrational bands which are observed in the FT-IR and FT-Raman spectra of the title compound are completely assigned for the first time with the help of TED. The donor–acceptor interaction, as obtained from NBO analysis could fairly explains the decrease of occupancies of σ bonding orbital and the increase of occupancy of π^* antibonding orbitals. The bioactivity of the molecule is proposed by means of band gap (−5.288 eV) energy derived from HOMO and LUMO calculation. The atomic charges of the present molecule has been calculated and also plotted.

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